



Smoothing the Landscape:

Moreau–Yosida Regularization in Modern
Density-Functional Theory

EPFL, September 16th, 2025

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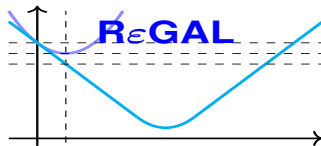
Acknowledgements

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Funded under ERC StG No. 101041487 REGAL



European Research Council

Established by the European Commission

Outline

- 1 Introduction to density-functional theory
- 2 Moreau–Yosida regularisation in DFT
- 3 Moreau–Yosida inversion scheme
- 4 References

Introduction to density-functional theory

$$\hat{H}(v) = \overbrace{-\frac{1}{2} \sum_j \nabla_j^2 + \frac{1}{2} \sum_{k \neq j} \frac{1}{|\mathbf{r}_j - \mathbf{r}_k|}}^{\hat{H}_0} + \sum_j v(\mathbf{r}_j) \quad (\mathbf{r}_1, \dots, \mathbf{r}_N) \in \mathbb{R}^{3N}$$

$$\psi : \mathbb{R}^{3N} \rightarrow \mathbb{C}, \quad \rho_\psi = N \int_{\mathbb{R}^{3(N-1)}} |\psi|^2 \quad \langle v, \rho \rangle = \int v \rho$$

$$\begin{aligned} E(v) &= \inf_{\psi} \langle \psi | \hat{H}(v) | \psi \rangle \\ &= \inf_{\rho} \{ F(\rho) + \langle v, \rho \rangle \} \end{aligned}$$

$$\begin{aligned} F(\rho) &= \sup_v \{ E(v) - \langle v, \rho \rangle \} \\ &= \inf_{\Gamma \mapsto \rho} \text{Tr } \hat{H}_0 \Gamma \\ &= \inf_{\rho = \sum_j w_j \rho_{\psi_j}} \sum_j w_j \langle \psi_j | \hat{H}_0 | \psi_j \rangle \end{aligned}$$

Hohenberg–Kohn theorem (conjecture?)

Suppose ρ_{gs} is a ground density of $\hat{H}(v_i)$, $i = 1, 2$

$$E(v_i) = \underbrace{\min_{\Gamma \mapsto \rho_{\text{gs}}} \text{Tr } \hat{H}_0 \Gamma + \langle v_i, \rho_{\text{gs}} \rangle}_{\text{independent of } v_i}$$

$\implies \hat{H}(v_1)$ and $\hat{H}(v_2)$ share a ground state $\psi \implies$

$$\text{subtracting} \quad \begin{cases} \hat{H}(v_1)\psi &= E(v_1)\psi \\ \hat{H}(v_2)\psi &= E(v_2)\psi \end{cases} \implies$$

$$(v_1(\mathbf{r}_1) - v_2(\mathbf{r}_1))\psi = \underbrace{\left(E(v_1) - E(v_2) - \sum_{j=2}^N (v_1(\mathbf{r}_j) - v_2(\mathbf{r}_j)) \right)}_{\text{independent of } x_1} \psi$$

If $\psi \neq 0$ (a.e.), “UCP” in the math. litt., then $v_1 - v_2 = \text{constant}$.

DFT: (F, E) conjugate pair

$$F(\rho) = \sup_v \{E(v) - \langle v, \rho \rangle\}$$

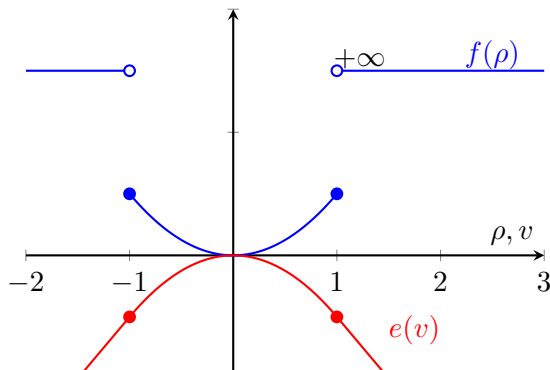
$$E(v) = \inf_{\rho} \{F(\rho) + \langle v, \rho \rangle\}$$

In general: $E(v) - F(\rho) \leq \langle v, \rho \rangle$

and equality for a ground-state pair (v, ρ)

Example: conjugation of a simple function

$$f(\rho) = \begin{cases} \frac{1}{2}\rho^2, & \text{if } |\rho| \leq 1, \\ \infty, & \text{if } |\rho| > 1 \end{cases} \Rightarrow e(v) = \inf_{\rho} (f(\rho) + v\rho) = \begin{cases} \frac{1}{2} - v, & \text{if } v > 1, \\ -\frac{1}{2}v^2, & \text{if } |v| \leq 1, \\ \frac{1}{2} + v, & \text{if } v < -1. \end{cases}$$



Moreau–Yosida regularisation in DFT

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Moreau–Yosida regularisation

Duality mapping $J_{\text{dual}} : X \rightarrow X^*$ (X str. convex and reflex., X^* uni. convex)

Density space: $N\text{-rep} \subsetneq X \ni \rho$

Potential space: $X^* \ni v$

Definition

Let $F : X \rightarrow \mathbb{R} \cup \{+\infty\}$ convex and lower semicontinuous functional. For some $\varepsilon > 0$, the *Moreau–Yosida regularisation* of F is

$$F^\varepsilon(\rho) = \inf_{\sigma \in X} \left\{ F(\sigma) + \frac{1}{2\varepsilon} \|\sigma - \rho\|_X^2 \right\}$$

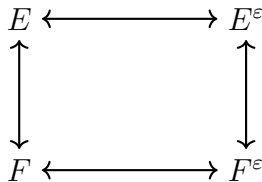
$$\rho^\varepsilon = \operatorname{argmin}_{\sigma \in X} \left\{ F(\sigma) + \frac{1}{2\varepsilon} \|\sigma - \rho\|_X^2 \right\}$$

Why? $F(\rho)$ in std DFT is **everywhere** discontinuous, P. Lammert [1]
Introduced into DFT by Kvaal et al. [2]

A Lossless Regularisation

$$E^\varepsilon(v) = \inf_{\rho \in X} \{F^\varepsilon(\rho) + \langle v, \rho \rangle\}$$

$$E^\varepsilon(v) = E(v) - \frac{\varepsilon}{2} \|v\|_V^2$$



The regularisation retains the ground-state energy and achieves diff.

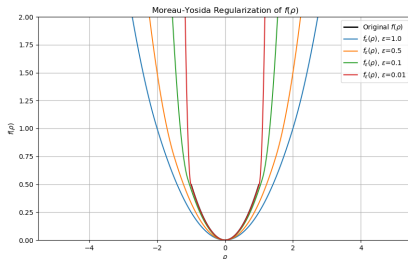
$$v^\varepsilon(\rho) = -\frac{\delta}{\delta \rho} F^\varepsilon(\rho) = \frac{1}{\varepsilon} J_{\text{dual}}[\rho^\varepsilon - \rho]$$

Quantitative HK:

$$\|v^\varepsilon(\rho_1) - v^\varepsilon(\rho_2)\|_{X^*} \leq \frac{C}{\varepsilon} \|\rho_1 - \rho_2\|_X$$

Example: Regularization of a simple function

$$f(\rho) = \begin{cases} \frac{1}{2}\rho^2, & \text{if } |\rho| \leq 1, \\ \infty, & \text{if } |\rho| > 1 \end{cases} \Rightarrow f^\varepsilon(\rho) = \begin{cases} \frac{1}{2} + \frac{1}{2\varepsilon}(1 - \rho)^2, & \text{if } \rho \geq 1 + \varepsilon, \\ \frac{\rho^2}{2(1+\varepsilon)}, & \text{if } |\rho| \leq 1 + \varepsilon, \\ \frac{1}{2} + \frac{1}{2\varepsilon}(1 + \rho)^2, & \text{if } \rho \leq -1 - \varepsilon. \end{cases}$$



The Kohn–Sham Approach

Interacting electrons

$$\hat{H} = -\frac{1}{2} \sum_j \nabla_j^2 + \frac{1}{2} \sum_{k \neq j} \frac{1}{|\mathbf{r}_j - \mathbf{r}_k|} + \sum_j v_{\text{ext}}(\mathbf{r}_j)$$

Non-interacting electrons (KS system)

$$\hat{H}_{\text{KS}} = -\frac{1}{2} \sum_j \nabla_j^2 + \sum_j [v_{\text{ext}}(\mathbf{r}_j) + v_{\text{H}}(\mathbf{r}_j) + v_{\text{xc}}(\mathbf{r}_j)]$$

ρ_{gs}

Exchange-correlation potential $v_{\text{xc}} \leftarrow$ unknown

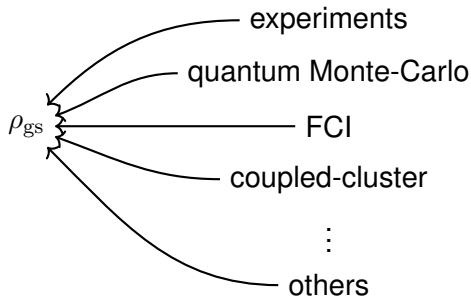
The forward problem:

Given a v_{ext} and an approximate E_{xc} , i.e., v_{xc} , determine ρ_{gs}

The Inverse Kohn–Sham Problem

The reverse problem:

Given a ρ_{gs} , determine v_{xc}



Why?

Better understanding of $F(\rho)$

Additional constraints for $F(\rho)$ to use
when designing approximations

Moreau–Yosida inversion scheme

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Kohn–Sham model functional

Recall

$$\hat{H}_{\text{KS}} = -\frac{1}{2} \sum_j \nabla_j^2 + \sum_j \overbrace{[v_{\text{ext}}(\mathbf{r}_j) + v_{\text{H}}(\mathbf{r}_j) + v_{\text{xc}}(\mathbf{r}_j)]}^{v_{\text{KS}}}$$

Given ρ_{gs} , let

$$\mathcal{F}(\rho) = T(\rho) + E_{\text{H}}(\rho) + \int v_{\text{ext}} \rho$$

Minimise

$$\mathcal{E}(\rho; \rho_{\text{gs}}) = \underbrace{\mathcal{F}(\rho)}_{\text{no xc}} + \underbrace{\frac{1}{2\varepsilon} \|\rho - \rho_{\text{gs}}\|_X^2}_{\text{targets xc as } \varepsilon \rightarrow 0+} \quad \text{over } \rho \in X \quad (1)$$

Proximal density, unique minimiser of $\mathcal{E}$:

$$\rho_{\text{gs}}^\varepsilon = \operatorname{argmin}_{\rho \in X} \mathcal{E}(\rho, \rho_{\text{gs}})$$

The Exchange-Correlation Potential

Since $\mathcal{F}^\varepsilon(\rho)$ is differentiable,

$$v_{\text{xc}}^\varepsilon = -\frac{\delta}{\delta\rho} \mathcal{F}^\varepsilon(\rho_{\text{gs}}) = \frac{1}{\varepsilon} J_{\text{dual}}[\rho_{\text{gs}}^\varepsilon - \rho_{\text{gs}}]$$

Theorem. [3, 4] Exchange-correlation potential obtained from

$$v_{\text{xc}}(\mathbf{r}) = \lim_{\varepsilon \rightarrow 0^+} v_{\text{xc}}^\varepsilon = \lim_{\varepsilon \rightarrow 0^+} \frac{1}{\varepsilon} J_{\text{dual}}[\rho_{\text{gs}}^\varepsilon - \rho_{\text{gs}}](\mathbf{r})$$

The form of J_{dual} depends on the choice of X :

Ex1. $X = L^2 \implies J_{\text{dual}} = I$ (van Leeuwen & Baerends)

Ex2. $X = \text{homogeneous Sobolev space} \implies J_{\text{dual}}(\rho) = v_{\text{H}}(\rho)$ (ZMP)

Sketch of proof: By assumption v_{KS} exists s.t. the KS system “gives” ρ_{gs} :

$$\min_{\rho} (T(\rho) + \langle v_{\text{KS}}, \rho \rangle) \iff \underline{\partial} T(\rho_{\text{gs}}) + v_{\text{KS}} \ni 0 \quad (2)$$

$$\min_{\rho} (\mathcal{F}(\rho) + \langle v_{\text{xc}}, \rho \rangle) \iff \underbrace{\underline{\partial} \mathcal{F}(\rho_{\text{gs}}) + v_{\text{xc}} \ni 0}_{(*)} \quad (3)$$

Moreau-Yosida:

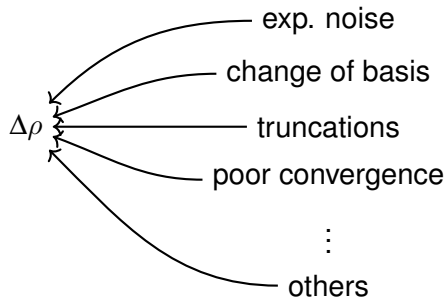
$$\min_{\rho} \left\{ \mathcal{F}(\rho) + \frac{1}{2\varepsilon} \|\rho - \rho_{\text{gs}}\|_{\mathcal{D}}^2 \right\} \iff \underbrace{\underline{\partial} \mathcal{F}(\rho_{\text{gs}}^{\varepsilon}) + \frac{1}{\varepsilon} J_{\text{dual}}(\rho_{\text{gs}}^{\varepsilon} - \rho_{\text{gs}}) \ni 0}_{(**)} \quad (4)$$

$$(*), (**) \text{ and } \rho_{\text{gs}}^{\varepsilon} \rightarrow \rho_{\text{gs}} \text{ as } \varepsilon \rightarrow 0+ \implies v_{\text{xc}} = \lim_{\varepsilon \rightarrow 0+} \frac{1}{\varepsilon} J_{\text{dual}}(\rho_{\text{gs}}^{\varepsilon} - \rho_{\text{gs}})$$

Error Bounds

Inexact density:

$$\tilde{\rho}_{\text{gs}} = \rho_{\text{gs}} + \Delta\rho$$



The *total error*

$$\|v_{\text{xc}} - \tilde{v}_{\text{xc}}^{\varepsilon}\|_{X^*} \leq \underbrace{\|v_{\text{xc}} - v_{\text{xc}}^{\varepsilon}\|_{X^*}}_{\text{terminating at finite } \varepsilon} + \underbrace{\|v_{\text{xc}}^{\varepsilon} - \tilde{v}_{\text{xc}}^{\varepsilon}\|_{X^*}}_{\text{using inexact } \rho_{\text{gs}}} \quad (5)$$

Energy error splitting

Kim, Sim and Burke [5]: Density-driven error

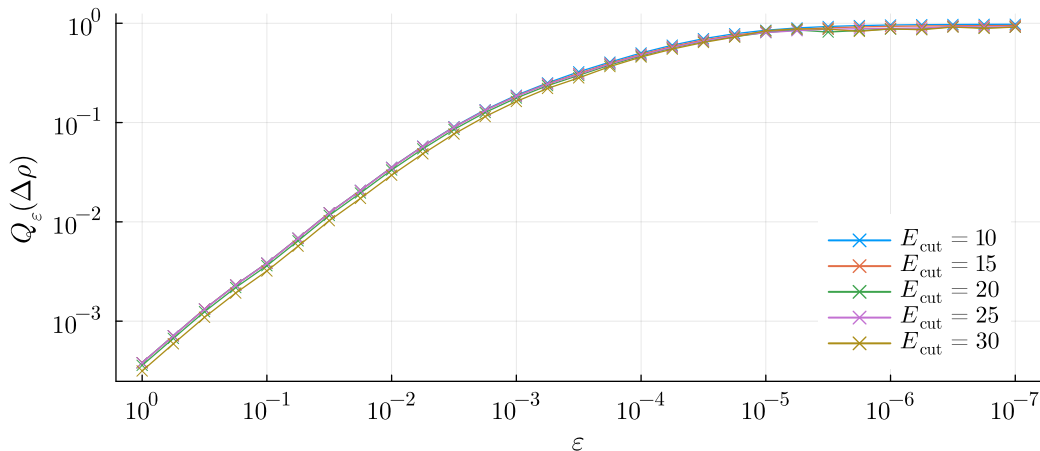
$$\begin{aligned}\Delta E &= E(\rho_{\text{gs}}) - E_{\text{appr}}(\rho_{\text{gs}} + \Delta\rho) \\ &= \underbrace{(E(\rho_{\text{gs}}) - E_{\text{appr}}(\rho_{\text{gs}}))}_{\text{error from functional}} + \underbrace{(E_{\text{appr}}(\rho_{\text{gs}}) - E_{\text{appr}}(\rho_{\text{gs}} + \Delta\rho))}_{\text{error from density}}\end{aligned}\quad (6)$$

We borrow the term “density-driven” when we speak of errors in the potentials caused by inexact densities:

$$\|v_{\text{xc}}^\varepsilon - \tilde{v}_{\text{xc}}^\varepsilon\|_{X^*} \leq \frac{1 + Q_\varepsilon}{\varepsilon} \|\Delta\rho\|_X$$

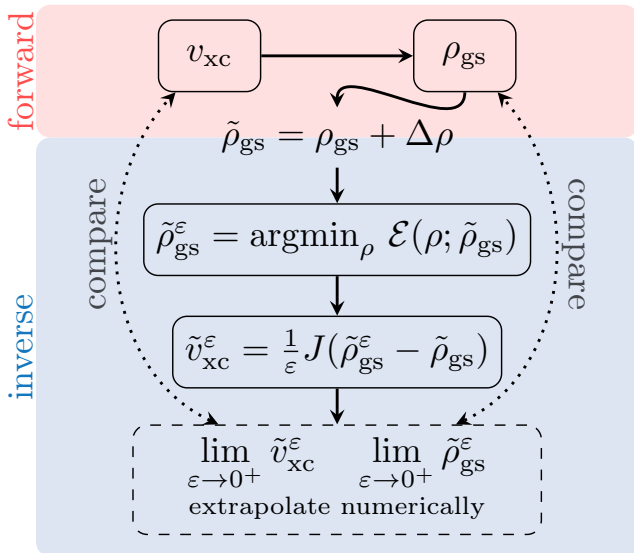
Non-Expansiveness of $\rho \mapsto \rho^\varepsilon$

$$\|\rho_{\text{gs}}^\varepsilon - \tilde{\rho}_{\text{gs}}^\varepsilon\|_X \leq \|\rho_{\text{gs}} - \tilde{\rho}_{\text{gs}}\|_X = \|\Delta\rho\|_X \quad Q_\varepsilon(\Delta\rho) = \frac{\|\rho_{\text{gs}}^\varepsilon - \tilde{\rho}_{\text{gs}}^\varepsilon\|_X}{\|\Delta\rho\|_X}$$

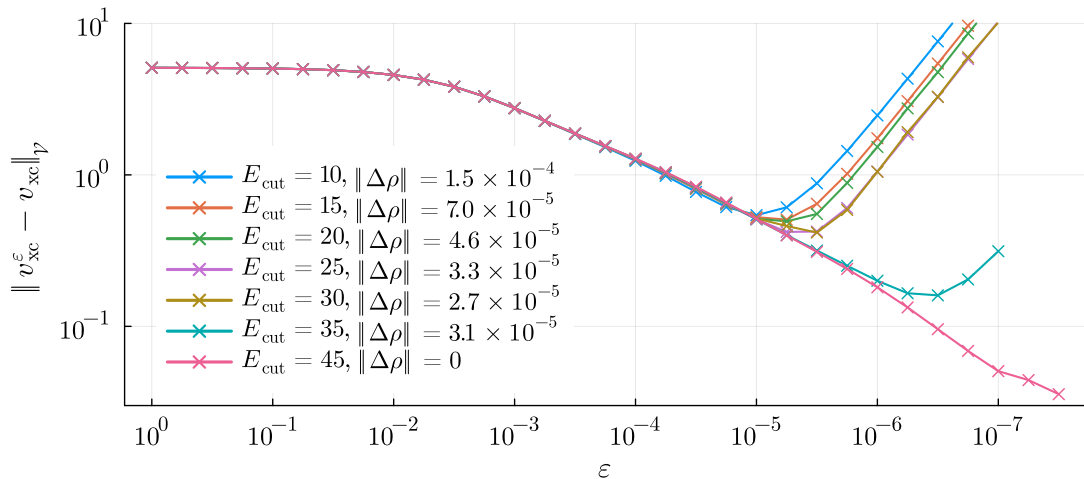


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The Inversion Algorithm with error $\Delta\rho$



Convergence: Bulk Silicon



Conclusion

- Regularization as a bridge between approximate and exact DFT
- Mathematical framework
- Algorithm for determining xc potentials
- Development of approximate functionals

References I

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4. **Herbst et al., M. F.** Kohn-Sham inversion with mathematical guarantees. *Phys. Rev. B* **111**, 205143 (2025).
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Thank you for your attention!

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